# Dynamics of Open Chains: Lagrangian Formulation

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# **Robot Dynamics**

- Study motion of robots with the forces and torques that cause them
- Equations of motion
  - A set of second-order differential equations

$$au = M( heta) \ddot{ heta} + h( heta, \dot{ heta})$$
 Joint variables  $heta \in \mathbb{R}^n$ 

Joint forces and torques  $\ au\in \mathbb{R}^n$   $M( heta)\in \mathbb{R}^{n imes n}$  a symmetric positive-definite mass matrix

 $h(\theta, \dot{\theta}) \in \mathbb{R}^n$ 

forces that lump together centripetal, Coriolis, gravity, and friction terms that depend on  $\theta$  and  $\dot{\theta}$ 

# Forward and Inverse Dynamics

- Forward dynamics
  - Given robot state  $(\theta, \dot{\theta})$  and the joint forces and torques
  - Determine the robot's acceleration

$$\ddot{\theta} = M^{-1}(\theta) \left( \tau - h(\theta, \dot{\theta}) \right)$$

Simulation

- Inverse dynamics
  - Given robot state  $( heta, \dot{ heta})$  and a desired acceleration
  - Find the joint forces and torques

$$\tau = M(\theta)\ddot{\theta} + h(\theta, \dot{\theta})$$

# Robot Dynamics

- Lagrangian formulation
  - Kinetic energy and potential energy
- Newton-Euler formulation
  - F = ma

- Generalized coordinates
  - Choose a set of independent coordinates  $\,q \in \mathbb{R}^n\,$  that describes the system's configuration
- Generalized forces  $f \in \mathbb{R}^n$ 
  - Power  $f^{\mathrm{T}}\dot{q}$
- Lagrangian function  $\mathcal{L}(q,\dot{q}) = \mathcal{K}(q,\dot{q}) \mathcal{P}(q)$

Kinetic energy Potential energy

• Equations of motion

$$f = \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{q}} - \frac{\partial \mathcal{L}}{\partial q}$$

Euler-Lagrange equations with external forces

• Example: a particle of mass m constrained to move on a vertical line

Generalized coordinate: height of the particle  $x \in \mathbb{R}$ 

external force Newton's second law  $f - \mathfrak{m}g = \mathfrak{m}\ddot{x}$ Lagrangian formulation  $f(x,\dot{x}) = \mathcal{K}(x,\dot{x}) - \mathcal{D}(x) = \frac{1}{2}\mathfrak{m}\dot{x}^2$ 

$$\mathcal{L}(x,\dot{x}) = \mathcal{K}(x,\dot{x}) - \mathcal{P}(x) = \frac{1}{2}\mathfrak{m}\dot{x}^2 - \mathfrak{m}gx$$

$$f = \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{x}} - \frac{\partial \mathcal{L}}{\partial x} = \mathfrak{m} \ddot{x} + \mathfrak{m} g$$

 $\mathfrak{m}q$ 



A 2R open chain under gravity

#### Two links with point masses

$$\begin{bmatrix} x_1 \\ y_1 \end{bmatrix} = \begin{bmatrix} L_1 \cos \theta_1 \\ L_1 \sin \theta_1 \end{bmatrix},$$
  
$$\begin{bmatrix} \dot{x}_1 \\ \dot{y}_1 \end{bmatrix} = \begin{bmatrix} -L_1 \sin \theta_1 \\ L_1 \cos \theta_1 \end{bmatrix} \dot{\theta}_1$$
  
$$\begin{bmatrix} x_2 \\ y_2 \end{bmatrix} = \begin{bmatrix} L_1 \cos \theta_1 + L_2 \cos(\theta_1 + \theta_2) \\ L_1 \sin \theta_1 + L_2 \sin(\theta_1 + \theta_2) \end{bmatrix},$$
  
$$\begin{bmatrix} \dot{x}_2 \\ \dot{y}_2 \end{bmatrix} = \begin{bmatrix} -L_1 \sin \theta_1 - L_2 \sin(\theta_1 + \theta_2) \\ L_1 \cos \theta_1 + L_2 \cos(\theta_1 + \theta_2) \end{bmatrix} - L_2 \sin(\theta_1 + \theta_2) \\ L_2 \cos(\theta_1 + \theta_2) \end{bmatrix} \begin{bmatrix} \dot{\theta}_1 \\ \dot{\theta}_2 \end{bmatrix}$$

Generalized coordinate  $\theta = (\theta_1, \theta_2)$ Generalized forces  $\tau = (\tau_1, \tau_2)$  Joint torques

Lagrangian 
$$\mathcal{L}(\theta, \dot{\theta}) = \sum_{i=1}^{2} (\mathcal{K}_{i} - \mathcal{P}_{i})$$



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Lagrangian 
$$\mathcal{L}(\theta, \dot{\theta}) = \sum_{i=1}^{2} (\mathcal{K}_{i} - \mathcal{P}_{i})$$
  
 $\mathcal{K}_{1} = \frac{1}{2} \mathfrak{m}_{1}(\dot{x}_{1}^{2} + \dot{y}_{1}^{2}) = \frac{1}{2} \mathfrak{m}_{1}L_{1}^{2}\dot{\theta}_{1}^{2}$   
 $\mathcal{K}_{2} = \frac{1}{2} \mathfrak{m}_{2}(\dot{x}_{2}^{2} + \dot{y}_{2}^{2})$   
 $= \frac{1}{2} \mathfrak{m}_{2} \left( (L_{1}^{2} + 2L_{1}L_{2}\cos\theta_{2} + L_{2}^{2})\dot{\theta}_{1}^{2} + 2(L_{2}^{2} + L_{1}L_{2}\cos\theta_{2})\dot{\theta}_{1}\dot{\theta}_{2} + L_{2}^{2}\dot{\theta}_{2}^{2} \right)$   
 $\mathcal{P}_{1} = \mathfrak{m}_{1}gy_{1} = \mathfrak{m}_{1}gL_{1}\sin\theta_{1},$   
 $\mathcal{P}_{2} = \mathfrak{m}_{2}gy_{2} = \mathfrak{m}_{2}g(L_{1}\sin\theta_{1} + L_{2}\sin(\theta_{1} + \theta_{2}))$   
 $\tau_{i} = \frac{d}{dt}\frac{\partial\mathcal{L}}{\partial\dot{\theta}_{i}} - \frac{\partial\mathcal{L}}{\partial\theta_{i}}, \qquad i = 1, 2.$   
 $\tau_{1} = (\mathfrak{m}_{1}L_{1}^{2} + \mathfrak{m}_{2}(L_{1}^{2} + 2L_{1}L_{2}\cos\theta_{2} + L_{2}^{2}))\ddot{\theta}_{1}$   
 $+ \mathfrak{m}_{2}(L_{1}L_{2}\cos\theta_{2} + L_{2}^{2})\ddot{\theta}_{2} - \mathfrak{m}_{2}L_{1}L_{2}\sin\theta_{2}(2\dot{\theta}_{1}\dot{\theta}_{2} + \dot{\theta}_{2}^{2})$   
 $+ (\mathfrak{m}_{1} + \mathfrak{m}_{2})L_{1}g\cos\theta_{1} + \mathfrak{m}_{2}gL_{2}\cos(\theta_{1} + \theta_{2}),$   
 $\tau_{2} = \mathfrak{m}_{2}(L_{1}L_{2}\cos\theta_{2} + L_{2}^{2})\ddot{\theta}_{1} + \mathfrak{m}_{2}L_{2}^{2}\ddot{\theta}_{2} + \mathfrak{m}_{2}L_{1}L_{2}\dot{\theta}_{1}^{2}\sin\theta_{2}$   
 $+ \mathfrak{m}_{2}gL_{2}\cos(\theta_{1} + \theta_{2}).$ 

g



Product-  
velocity Gravity  
term term 
$$\tau = M(\theta)\ddot{\theta} + \underbrace{c(\theta, \dot{\theta}) + g(\theta)}_{h(\theta, \dot{\theta})},$$

$$M(\theta) = \begin{bmatrix} \mathfrak{m}_1 L_1^2 + \mathfrak{m}_2 (L_1^2 + 2L_1 L_2 \cos \theta_2 + L_2^2) & \mathfrak{m}_2 (L_1 L_2 \cos \theta_2 + L_2^2) \\ \mathfrak{m}_2 (L_1 L_2 \cos \theta_2 + L_2^2) & \mathfrak{m}_2 L_2^2 \end{bmatrix}$$

$$c(\theta, \dot{\theta}) = \begin{bmatrix} -\mathfrak{m}_2 L_1 L_2 \sin \theta_2 (2\dot{\theta}_1 \dot{\theta}_2 + \dot{\theta}_2^2) \\ \mathfrak{m}_2 L_1 L_2 \dot{\theta}_1^2 \sin \theta_2 \end{bmatrix}$$
 Coriolis and centripetal torques

A 2R open chain under gravity

$$g(\theta) = \begin{bmatrix} (\mathfrak{m}_1 + \mathfrak{m}_2)L_1g\cos\theta_1 + \mathfrak{m}_2gL_2\cos(\theta_1 + \theta_2) \\ \mathfrak{m}_2gL_2\cos(\theta_1 + \theta_2) \end{bmatrix} \text{ gravitational torques}$$



- For general n-link open chain and all joints are actuated
- Generalized coordinates are joint values  $\, \theta \, \in \, \mathbb{R}^n \,$
- Generalized forces  $\ \tau \in \mathbb{R}^n$ 
  - Revolute joint -> torque
  - Prismatic joint -> force

• Lagrangian 
$$\mathcal{L}(\theta, \dot{\theta}) = \mathcal{K}(\theta, \dot{\theta}) - \mathcal{P}(\theta)$$

$$\mathcal{K}(\theta) = \frac{1}{2} \sum_{i=1}^{n} \sum_{j=1}^{n} m_{ij}(\theta) \dot{\theta}_{i} \dot{\theta}_{j} = \frac{1}{2} \dot{\theta}^{\mathrm{T}} M(\theta) \dot{\theta}$$
$$n \times n \text{ mass matrix } M(\theta)$$

• Dynamic equations 
$$\tau_{i} = \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{\theta}_{i}} - \frac{\partial \mathcal{L}}{\partial \theta_{i}}, \qquad i = 1, \dots, n.$$

$$\tau_{i} = \sum_{j=1}^{n} m_{ij}(\theta) \ddot{\theta}_{j} + \sum_{j=1}^{n} \sum_{k=1}^{n} \Gamma_{ijk}(\theta) \dot{\theta}_{j} \dot{\theta}_{k} + \frac{\partial \mathcal{P}}{\partial \theta_{i}}, \qquad i = 1, \dots, n,$$
Christoffel symbols of the first kind 
$$\Gamma_{ijk}(\theta) = \frac{1}{2} \left( \frac{\partial m_{ij}}{\partial \theta_{k}} + \frac{\partial m_{ik}}{\partial \theta_{j}} - \frac{\partial m_{jk}}{\partial \theta_{i}} \right)$$

$$\tau = M(\theta) \ddot{\theta} + \dot{\theta}^{T} \Gamma(\theta) \dot{\theta} + g(\theta) \qquad \Gamma(\theta) \qquad n \times n \times n$$

$$g(\theta) = \frac{\partial \mathcal{P}}{\partial \theta}$$

$$\dot{\theta}^{T} \Gamma_{0}(\theta) \dot{\theta} = \begin{bmatrix} \dot{\theta}^{T} \Gamma_{1}(\theta) \dot{\theta} \\ \dot{\theta}^{T} \Gamma_{2}(\theta) \dot{\theta} \\ \vdots \\ \dot{\theta}^{T} \Gamma_{n}(\theta) \dot{\theta} \end{bmatrix} \qquad \Gamma_{i}(\theta) \text{ is an } n \times n \text{ matrix with } (j, k) \text{ th entry } \Gamma_{ijk}$$

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$$\tau = M(\theta)\ddot{\theta} + C(\theta, \dot{\theta})\dot{\theta} + g(\theta)$$

Coriolis matrix  $C( heta, \dot{ heta}) \in \mathbb{R}^{n imes n}$ 

$$c_{ij}(\theta, \dot{\theta}) = \sum_{k=1}^{n} \Gamma_{ijk}(\theta) \dot{\theta}_k$$

# Understanding the Mass Matrix

• Kinetic energy 
$$\mathcal{K}(\theta) = \frac{1}{2} \sum_{i=1}^{n} \sum_{j=1}^{n} m_{ij}(\theta) \dot{\theta}_i \dot{\theta}_j = \frac{1}{2} \dot{\theta}^{\mathrm{T}} M(\theta) \dot{\theta}_i$$

- A generalization of point mass  $\frac{1}{2}mv^{\mathrm{T}}v$
- M is positive definite  $\dot{\theta}^{\mathrm{T}} M(\theta) \dot{\theta} > 0$  for all  $\dot{\theta} \neq 0$ 
  - Mass of a point mass is always positive

# Understanding the Mass Matrix

- For a point mass  $f=m\ddot{x}$ 
  - m is independent of the acceleration direction
- A mass matrix presents different effective mass in different acceleration directions

$$\tau = M(\theta)\ddot{\theta} + C(\theta, \dot{\theta})\dot{\theta} + g(\theta)$$

### Understanding the Mass Matrix

 $\dot{\theta} = 0$ 



# Summary

- Robot Dynamics
- Lagrangian formulation

# Further Reading

• Chapter 8 in Kevin M. Lynch and Frank C. Park. Modern Robotics: Mechanics, Planning, and Control. 1st Edition, 2017.